



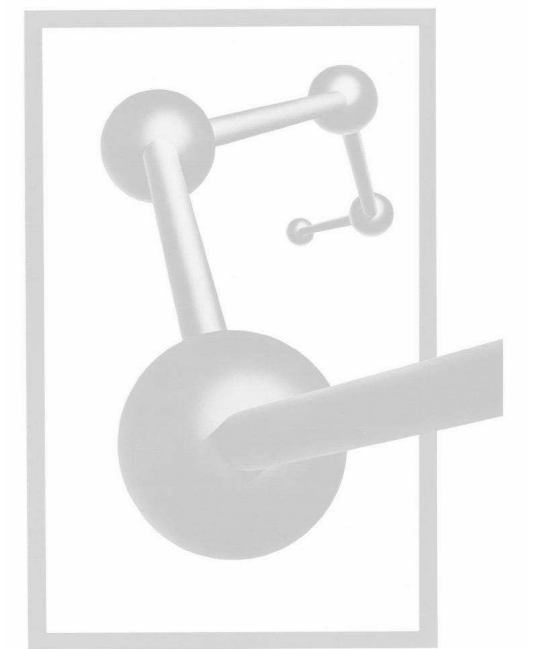
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# Exotic Ion $H_3^{++}$ in Strong Magnetic Fields: low-lying states

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## Abstract

### Main Results:

- (1) The exotic system  $H_3^{++}$  (which does not exist without magnetic field) exists in strong magnetic fields: (a) in triangular configuration for  $B \approx 10^8 - 10^{11}$  G and (b) in linear configuration for  $B > 10^{10}$  G
- (2) In the linear configuration the positive  $z$ -parity states  $1\sigma_g, 1\pi_u, 1\delta_g$  are bound states.
- (3) In the linear configuration the negative  $z$ -parity states  $1\sigma_u, 1\pi_g, 1\delta_u$  are repulsive states.

### Possible application:

content of a neutron star atmosphere (about  $\sim 10$  Km in radius and  $\sim 10$  cm of atmosphere thickness) under a strong surface magnetic field  $B = 10^{12} - 10^{13}$  G.

In 70's the first theoretical arguments were given indicating the physics of atoms and molecules in a strong magnetic field can exhibit a wealth of new and unexpected phenomena (Kadomtsev-Kudryatsev-Ruderman). In particular, the possibility of formation of unusual chemical compounds not existing without a magnetic field was emphasized. In practice, the atmosphere of neutron stars, which is characterized by enormous surface magnetic fields  $\sim 10^{12} - 10^{13}$  G (up to  $10^{15}$  G for magnetars), provides a valuable paradigm where this physics can be realized.

Recently the observational data (collected by the Chandra X-ray spatial observatory) of the soft X-ray spectrum of the isolated neutron star 1E1027.4-5209 revealed for the first time certain irregularities (Sanwal et al. '02). These irregularities might be interpreted as absorption features at  $\sim 0.7$  KeV and  $\sim 1.4$  KeV of a possible atomic-molecular nature assigned to the content of the neutron star atmosphere. In this context the study of the atomic and molecular, traditional and exotic systems in strong magnetic fields deserves a special attention.

The first already observed features for some atomic-molecular systems when the magnetic field increases are:

- The *total* and *binding* energies increase (The systems become more and more bound)
- Drastic decrease in the electron localization length in both directions, transverse and longitudinal, leading to a decrease in the equilibrium distance (The systems become more and more compact)
- The electronic cloud takes a needle-like form forming in most cases the quasi-one-dimensional systems with very unusual physical properties.

## Goal

Study of molecular systems, traditional and exotic, in strong magnetic fields in the context of the variational method with simple and unique trial functions. These functions are modifications of the celebrated Heitler-London, Hund-Mulliken and Guillemin-Zener functions. Our goal is to find the domain of existence of the systems and then to obtain their description with sufficiently high accuracy in the whole range of magnetic fields where the non-relativistic approximation is valid:  $B = 0 - 4.414 \times 10^{13}$  G.

In particular, we have carried out a detailed study in the Born-Oppenheimer approximation of the ground state for the one-electron systems:  $(ppe)$ -system - the traditional molecular ion  $H_2^+$  (the simplest molecular system) and for the systems  $(pppe)$  and  $(ppppe)$  giving rise to the exotic molecular ions  $H_3^{++}$  and  $H_4^{+++}$ , respectively, in linear configuration parallel to the magnetic field (Fig.1), and for different spatial configurations of the charged centers.

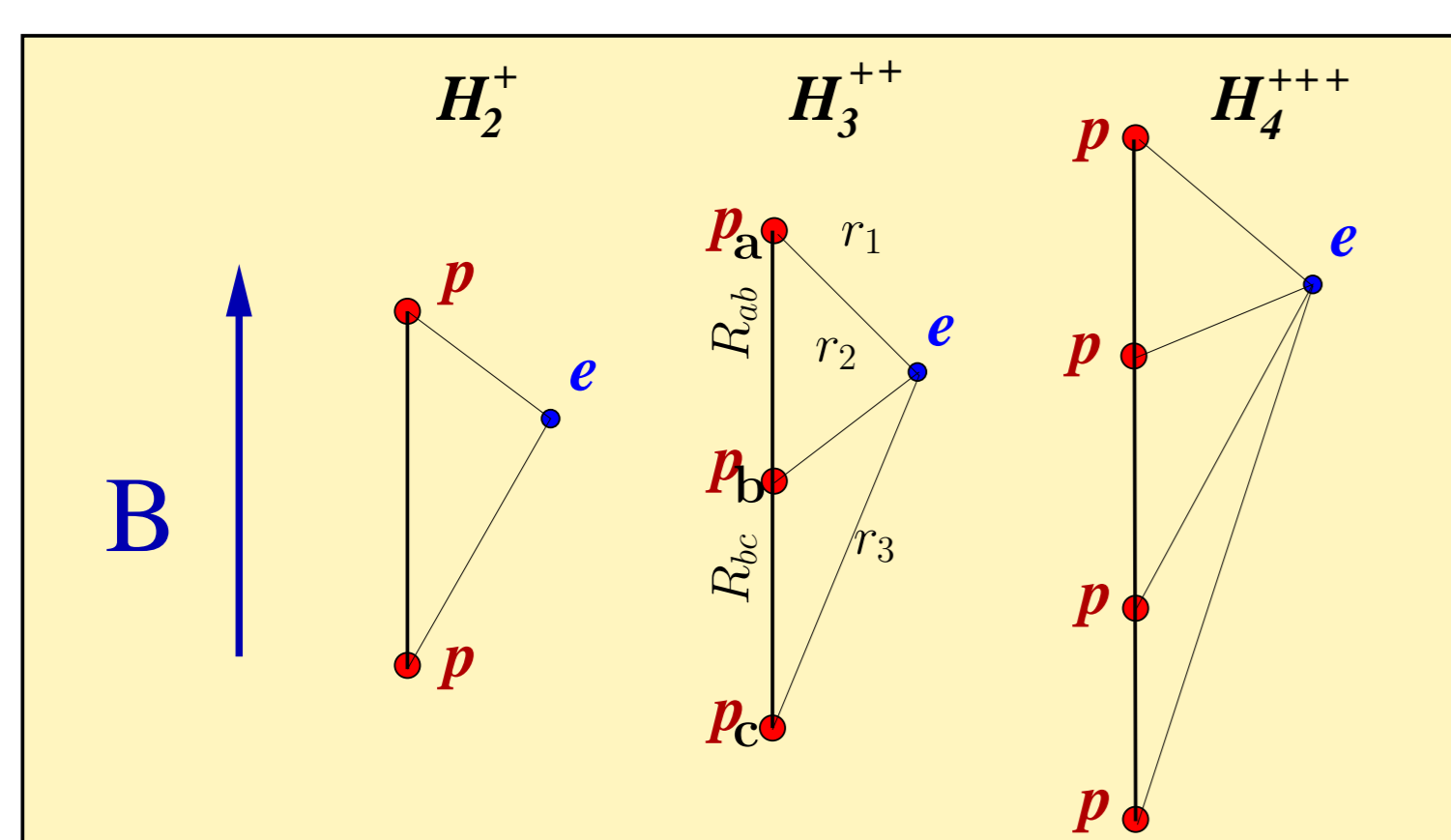


Fig. 1: One-electron molecular systems in a uniform magnetic field  $\mathbf{B} = (0, 0, B)$ : the traditional system  $H_2^+$ , and the exotic ions  $H_3^{++}$  and  $H_4^{+++}$ .

## Formulation of the Problem

In the Born-Oppenheimer approximation of zero order (where protons are infinitely massive), the Schrödinger equation for a system of one-electron and a certain number of protons in a magnetic field  $\mathbf{B}$  is

$$\left[ (\mathbf{p} - \mathbf{A})^2 + V_c(r_1, r_2, \dots) \right] \Psi(\mathbf{r}) = E\Psi(\mathbf{r}) \quad (\text{in a.u.})$$

where  $\mathbf{p} = -i\nabla$  is the *momentum of the electron*,  $\mathbf{A}$  is a *vector potential* associated with the constant uniform magnetic field  $\mathbf{B}$ , and  $V_c(r_1, r_2, \dots)$  describes the *Coulomb interaction* between the charged particles. For the  $(pppe)$ -system, (see Figs.1,2):

$$V_c = \frac{2}{R_{ab}} + \frac{2}{R_{ac}} + \frac{2}{R_{bc}} - \frac{2}{r_1} - \frac{2}{r_2} - \frac{2}{r_3}.$$

In particular, for a uniform magnetic field in the  $z$ -direction,  $\mathbf{B} = (0, 0, B)$  we can choose the vector potential in a *gauge*:

$$\mathbf{A} = B((\xi - 1)y, \xi x, 0),$$

where  $\xi$  is a gauge parameter. Observables are, of course, gauge invariant and do not depend on the value of  $\xi$ . In variational calculations for a fixed class of trial functions the variational energy (which is approximate) can depend on the gauge, there can exist an optimal value of  $\xi$  giving a minimal variational energy. Therefore one can consider  $\xi$  as a variational parameter.

### $H_3^{++}$ (triangular)

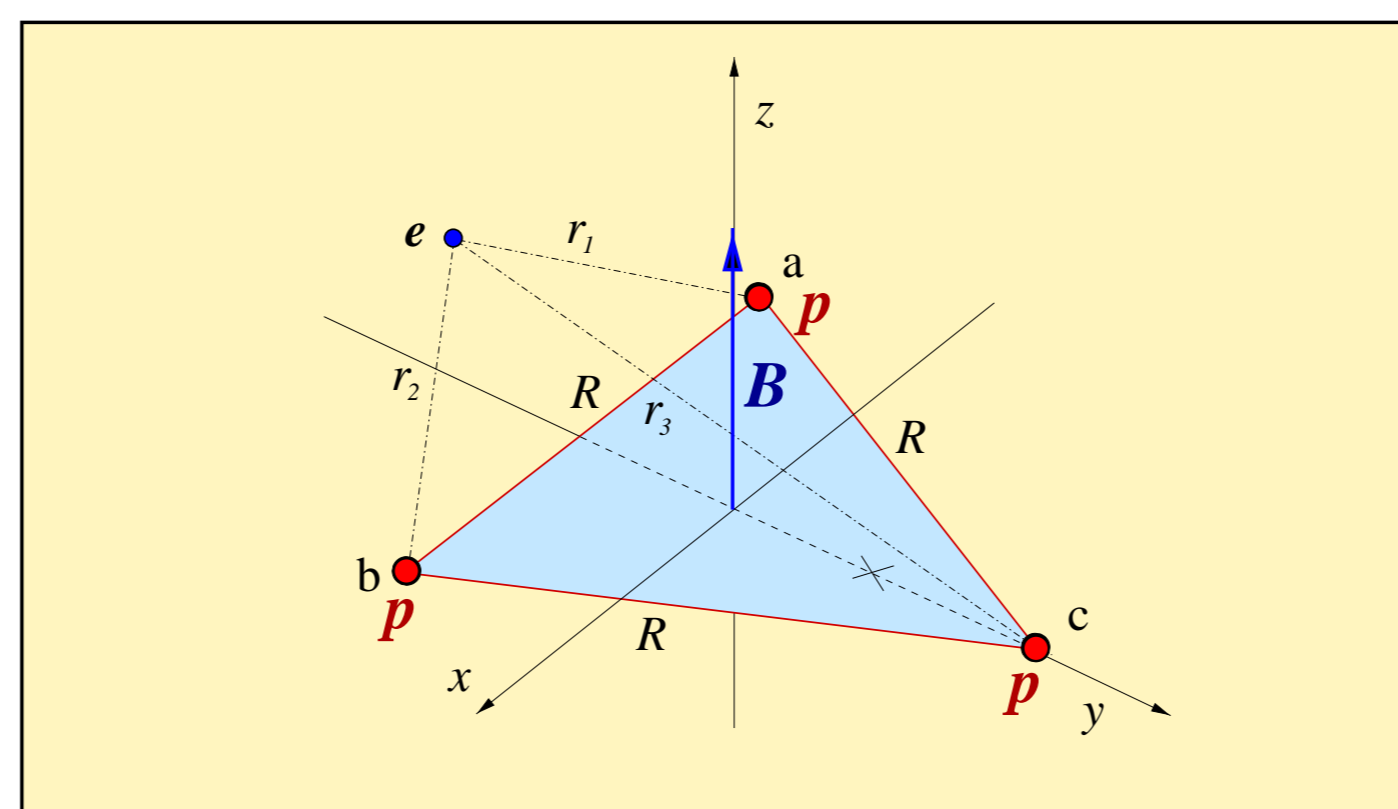


Fig.2: The  $H_3^{++}$  molecular ion in triangular (equilateral) configuration perpendicular to the magnetic field  $\mathbf{B}$ .

1. The system  $(pppe)$  in triangular equilateral configuration (See Fig.2) develops a well pronounced minimum for  $10^8 \text{ G} < B < 10^{11} \text{ G}$ . Therefore the system  $H_3^{++}$  can exist for those magnetic fields
2. As the magnetic field increases the binding energy increases and the size of the triangle diminishes.

### $H_3^{++}$ (linear)

1. The system  $H_3^{++}$  in linear configuration exists for  $B \geq 10^{11} \text{ G}$ .
2. As the magnetic field increases the binding energy increases and the internuclear distance decreases
3. For  $B > 10^{13} \text{ G}$ ,  $H_3^{++}$  is the most stable one-electron system made from protons, in particular,  $E_T^{H_3^{++}} < E_T^{H_2^+}$ .

E.g. the binding energies for  $B = 4.414 \times 10^{13} \text{ G}$  are

$$E_b^{H_3^{++}} = 55.23 \text{ Ry}, E_b^{H_2^+} = 54.50 \text{ Ry}, E_b^{H-atom} = 32.77 \text{ Ry}$$

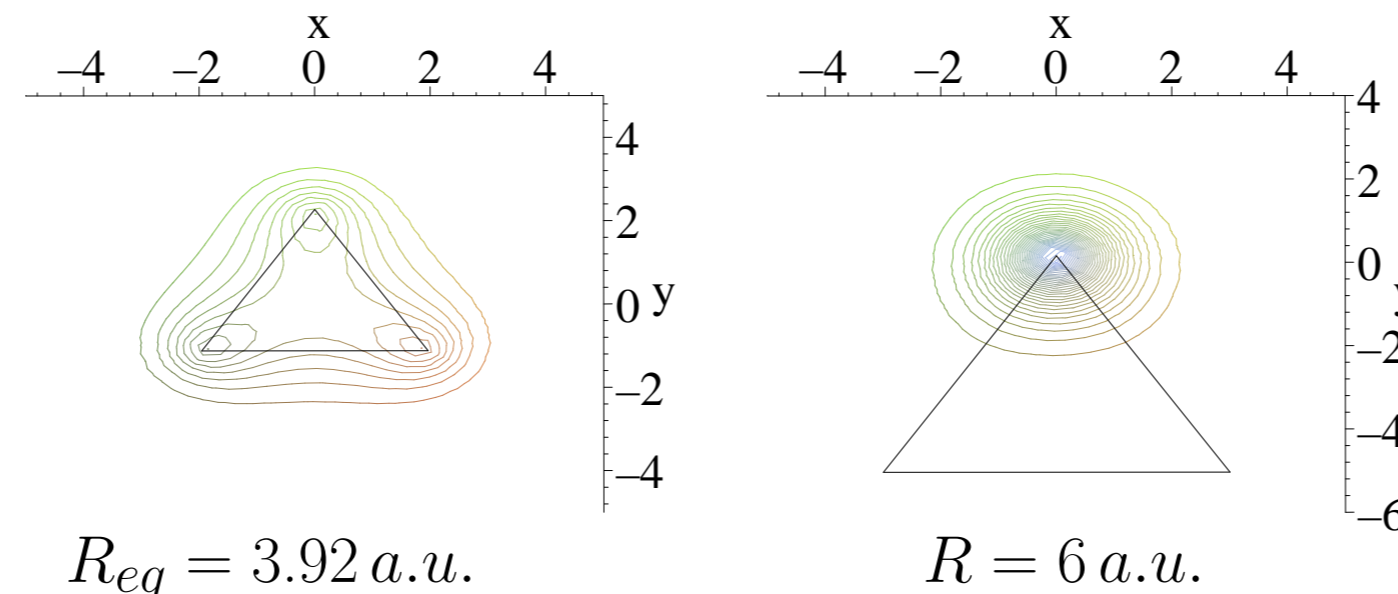


Fig. 3: Evolution of the electronic distribution as function of the size of the triangle (of side  $R$ ) in a magnetic field  $B = 10^9 \text{ G}$ . Plots on the left describe the bound state and the plots in the right correspond to the situation when the dissociation  $H_3^{++} \rightarrow H + p + p$  occurs.

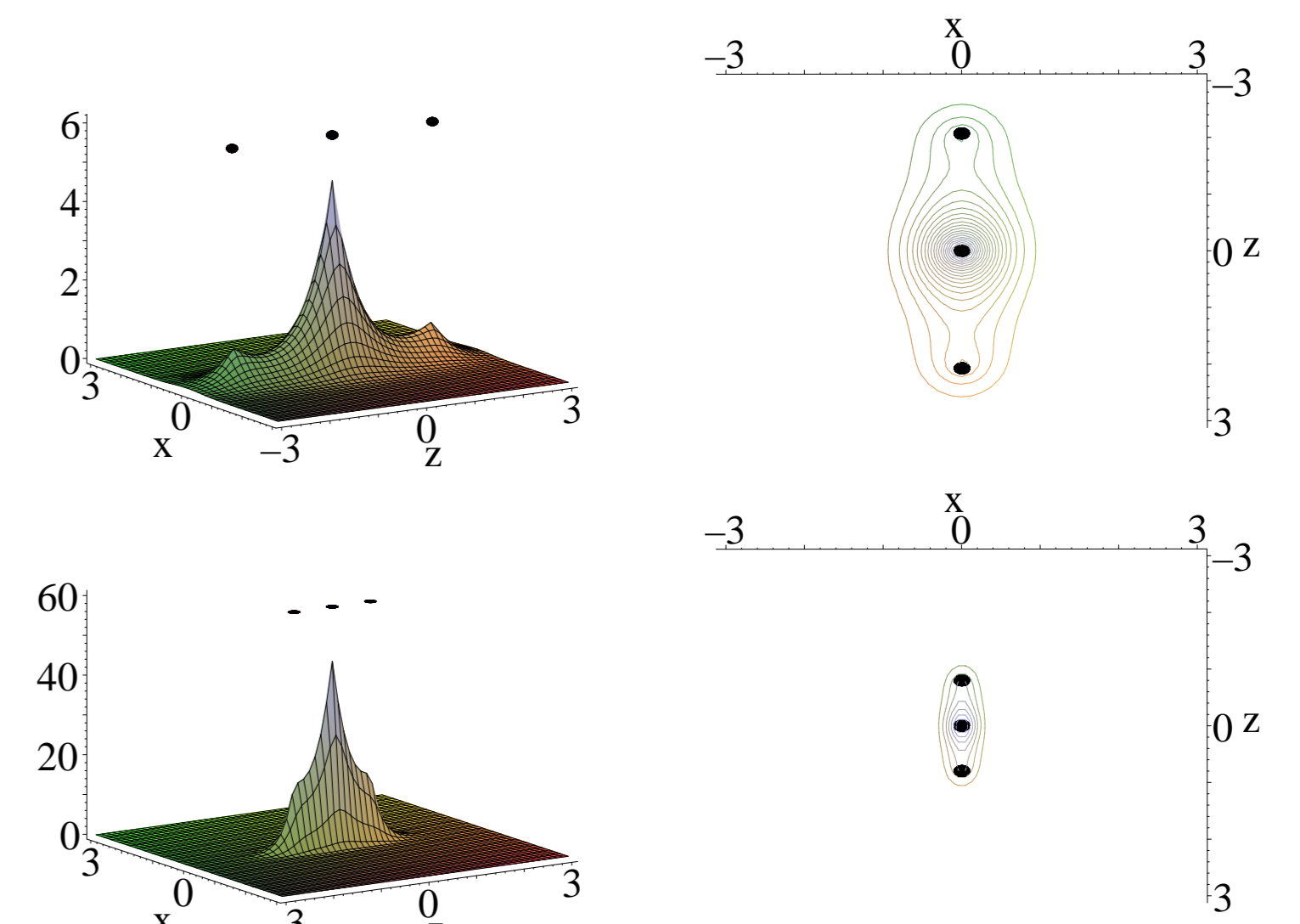


Fig. 4: Illustration of the evolution of the normalized electronic distribution  $\psi^2(x, y = 0, z) / \int \psi^2(x, y, z) d\mathbf{r}$  for the ion  $H_3^{++}$  in linear configuration aligned with a magnetic field at (a)  $B = 10^{11} \text{ G}$  and (b)  $B = 10^{12} \text{ G}$ .

## Excited states of $H_3^{++}$

For all magnetic fields  $B \geq 10^{11} \text{ G}$ :

1. Positive  $z$ -parity states  $1\sigma_g, 1\pi_u, 1\delta_g$  are bound
2. Negative  $z$ -parity states  $1\sigma_u, 1\pi_g, 1\delta_u$  are repulsive
3. The following hierarchy appears:

$$E_T^{1\sigma_g} < E_T^{1\pi_u} < E_T^{1\delta_g}, \quad R_{eq}^{1\sigma_g} < R_{eq}^{1\pi_u} < R_{eq}^{1\delta_g}.$$

E.g. for  $B = 10^{12} \text{ G}$ :

$$E_b^{1\sigma_g} = 15.16 \text{ Ry}, E_b^{1\pi_u} = 10.17 \text{ Ry}, E_b^{1\delta_g} = 8.28 \text{ Ry}$$

$$R_{eq}^{1\sigma_g} = 0.345 \text{ a.u.}, R_{eq}^{1\pi_u} = 0.497 \text{ a.u.}, R_{eq}^{1\delta_g} = 0.601 \text{ a.u.}$$

The absorption features at  $\sim 700 \text{ eV}$  and  $\sim 1400 \text{ eV}$  discovered in radiation of the neutron star 1E1207.4-5209 by Chandra (Sanwal et al., 2002) and confirmed by XMM-Newton can be explained if one assumes that the main abundance in the atmosphere is  $H_3^{++}$  and the magnetic field strength is  $(4 \pm 2) \times 10^{14} \text{ G}$ . Photodissociation  $H_3^{++} \rightarrow H + 2p$  corresponds to absorption feature at  $\sim 700 \text{ eV}$ , and photoionization  $H_3^{++} \rightarrow e + 3p$  contributes to absorption feature at  $\sim 1400 \text{ eV}$ . Photodissociation  $H_3^{++} \rightarrow H_2^+ + p$  corresponds to absorption feature at  $80 - 150 \text{ eV}$ , and it is not seen being beyond of Chandra detector acceptance.

## Publications

1. 'One-electron molecular systems in a strong magnetic field', A. Turbiner and J.C. López Vieyra *Physics Reports*, vol. 424, (March 2006) pp. 309-396
2. (a) ' $H_3^{++}$  can exist in strong magnetic field', A. Turbiner, J.-C. Lopez and U. Solis H., *JETP Letters* **69**, 844-850 (1999)  
(b) 'The exotic  $H_3^{++}$  ion in a strong magnetic field. Linear configuration', A.V. Turbiner, J.C. López Vieyra and N. L. Guevara, *Physical Review A* **72**, 023403, (2005)
3. ' $H_3^{++}$  in a strong magnetic field: triangular configuration', J.C. Lopez V. and A. Turbiner, *Physical Review A* **66**, 023409 (2002)